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(1940 - 2021)

Canonical Pairs

$$[q, p] = i\hbar \quad (\hbar = 2)$$

Stone-v. Neumann
Theorem

Unique $\sim (\hat{x}, -2i\partial)$ on $L^2(\mathbb{R}, dx)$

Schrödinger
Rep.

$$\begin{cases} \hat{x} f(x) = x f(x) \\ \partial f(x) = f'(x). \end{cases}$$

Creation / Annihilation

$$[a, a^*] = 1$$

$$q = a + a^*$$

$$p = \frac{1}{i}(a - a^*)$$

Number

$$N = a^* a \geq 0$$

$$\text{spectrum } N = \{0, 1, 2, 3, \dots\}.$$

$$\mathbb{E}[\psi(q, p)] \equiv \text{tr}[\rho \psi(\hat{x}, -2i\partial)] \text{ on } L^2(\mathbb{R})$$

In the Schrödinger rep.

$|n\rangle_{\text{number}}$ — n^{th} Hermite function

$$N|n\rangle_{\text{number}} = n|n\rangle_{\text{number}}$$

$$a^*|n\rangle_{\text{number}} = \sqrt{n+1}|n+1\rangle_{\text{number}}$$

$$a|n\rangle_{\text{number}} = \begin{cases} \sqrt{n}|n-1\rangle_{\text{number}} & (n \geq 1); \\ 0 & n = 0. \end{cases}$$

Exponential vectors ($\alpha \in \mathbb{C}$)

$$|\alpha\rangle = \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} |n\rangle_{\text{number}},$$

$$\langle \alpha | \beta \rangle = e^{\alpha^* \beta}$$

$$a|\alpha\rangle = \alpha|\alpha\rangle$$

In the ground state $|0\rangle = |0\rangle_{\text{number}}$

* q, p are Gaussian ($\mu=0, \sigma=1$).

In the state $\mathbb{E}_\alpha[\cdot] = \frac{\langle \alpha | \cdot | \alpha \rangle}{\langle \alpha | \alpha \rangle}$

* q, p are Gaussian ($\mu = 2\text{Re}/\text{Im} \alpha, \sigma=1$)

* N is Poissonian (intensity $|\alpha|^2$)

$$e^{i\theta N} |\alpha\rangle = \sum_n e^{i\theta n} \frac{\alpha^n}{\sqrt{n!}} |n\rangle_{\text{number}}$$

$$= |e^{i\theta} \alpha\rangle,$$

$$\text{so } \mathbb{E}[e^{i\theta N}] = \frac{e^{\alpha^* (e^{i\theta} \alpha)}}{e^{\alpha^* \alpha}} = e^{|\alpha|^2 (e^{i\theta} - 1)}.$$

! The most important prob. laws (continuous + discrete) are already there in undergraduate quantum physics!

Wigner's Quasi-probability: \mathcal{W}_ψ

$$\langle \psi | e^{iuq + ivp} | \psi \rangle$$

$$\equiv \int_{\mathbb{R}^2} e^{iux + ivy} \mathcal{W}_\psi(x, y) dx dy$$

Hudson's Theorem (1974)

$\mathcal{W}_\psi \geq 0 \Leftrightarrow \psi$ is Gaussian.

$$\text{i.e. } \psi(x) = e^{-ax^2 + bx + c}$$

for some $a, b, c \in \mathbb{C}$,
 $\operatorname{Re} a > 0$.

Quantum Central Limit Theorem

Let $(q_k, p_k)_{k=1}^{\infty}$ be an i.i.d. sequence of canonical pairs.

Independent?

$$\mathbb{E}[f(q_1, p_1, q_2, p_2)] = \int_{\rho_{12}} f(x_1, -2i\partial_1, x_2, -2i\partial_2)$$

then $\rho_{12} = \rho_1 \otimes \rho_2$ on $\underline{L^2(\mathbb{R}, dx_1)} \otimes \underline{L^2(\mathbb{R}, dx_2)}$

$$\text{Let } Q_n = \frac{1}{\sqrt{n}} \sum_{k=1}^n q_k, P_n = \frac{1}{\sqrt{n}} \sum_{k=1}^n p_k$$

$$\text{then } [Q_n, P_n] = 2i.$$

$$\text{Suppose } \mathbb{E}[q_n] = \mathbb{E}[p_n] = 0,$$

$$\mathbb{E}[q_n^2] = \mathbb{E}[p_n^2] = \sigma^2, \quad \mathbb{E}[q_n p_n + p_n q_n] = 0.$$

Theorem (Cushner + Hudson, 1971)

$$\lim_{n \rightarrow \infty} \mathbb{E} [e^{iuQ_n + ivP_n}] = \tilde{\mathbb{E}} [e^{iu\tilde{Q} + iv\tilde{P}}]$$

$$* [\tilde{Q}, \tilde{P}] = 1$$

* $\sigma = 1$, take Schrod. rep on $L^2(\mathbb{R})$
 $\tilde{\mathbb{E}}$ - groundstate $|0\rangle$

* $\sigma > 1$, take $L^2(\mathbb{R}) \otimes L^2(\mathbb{R})$

$$\tilde{Q} = \frac{\sqrt{\sigma^2 + 1}}{2} q \otimes 1 + \frac{\sqrt{\sigma^2 - 1}}{2} 1 \otimes q$$

$$\tilde{P} = \frac{\sqrt{\sigma^2 + 1}}{2} p \otimes 1 + \frac{\sqrt{\sigma^2 - 1}}{2} 1 \otimes p$$

and $\tilde{\mathbb{E}}$ - $|0\rangle \otimes |0\rangle$.

N.B. $\sigma^2 = \Delta q \Delta p \geq 1$ (HUP).

Quantum Ito Calculus

(Fock)
Fockspace

Bosons!

$$\Gamma(h) = (\mathbb{C}|\Omega\rangle) \oplus h \oplus \left(\bigoplus_{\text{sym.}}^2 h \right) \oplus \dots$$

$$|f\rangle = \Omega \oplus f \oplus \left(\frac{1}{\sqrt{2}} f \otimes f \right) \oplus \dots$$

* Annihilator $A(g)|f\rangle = \langle g|f\rangle |f\rangle$

* Creator $A(g)^*|f\rangle = \left. \frac{\partial}{\partial u} |f + ug\rangle \right|_{u=0}$

* Conserved /
Gauge /
Number $\Lambda(H)|f\rangle = \left. \frac{1}{i} \frac{\partial}{\partial u} |e^{iH_u} f\rangle \right|_{u=0}$

$$[A(t), A^*(q)] = \langle \delta | \delta \rangle$$

$$\text{Let } h = L^2(\mathbb{R}_+, dt)$$

$$A_t = A(\mathbb{1}_{[0, t]})$$

$$A_t^* = A^*(\mathbb{1}_{[0, t]})$$

$$\Lambda_t = \Lambda(\mathbb{I}_{[0, t]})$$

Quantum Itô Table

$$dA_t dA_t^* = dt$$

$$dA_t d\Lambda_t = dA_t$$

$$d\Lambda_t dA_t^* = dA_t^*$$

$$d\Lambda_t d\Lambda_t = d\Lambda_t$$

Hudson + Parthasarathy (1984)

on $h \otimes \Gamma(L^2(\mathbb{R}_+, dA_t))$

$$dU_t = \left[(S-1) \otimes dA_t + L \otimes dA_t^* - LS^* \otimes dA_t - \left(\frac{1}{2}L^*L + iH\right) \otimes dt \right] U_t$$

$$U_0 = I \otimes I .$$

(S, L, H) in $\mathcal{B}(h)$

$$S^{-1} = S^* , H = H^* .$$

$(U_t)_{t \geq 0}$ extends to a unitary

$$\langle u \otimes \Omega | U_t^*(X \otimes I) U_t | v \otimes \Omega \rangle$$

$$\equiv \langle u | \Phi_t(X) v \rangle , \forall u, v \in h$$

Φ_t a $\mathcal{C}P$ semigroup with generator

$$\mathcal{L}(X) = \frac{1}{2} [L^*, X]L + \frac{1}{2} L^* [X, L] - i [X, H] .$$

Applications

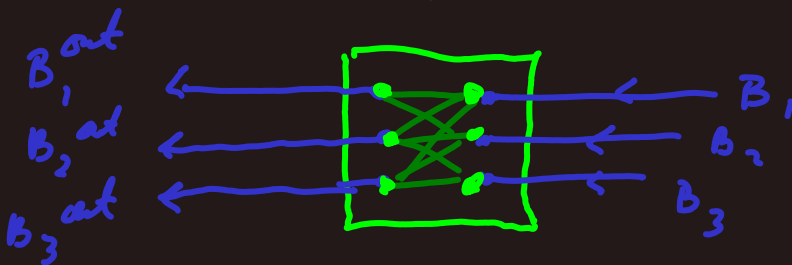
$$h \otimes \Gamma(\mathbb{C}^n \otimes L^2(\mathbb{R}_+))$$

$$B(t) = \begin{bmatrix} B_1(t) \\ \vdots \\ B_n(t) \end{bmatrix}, \quad \Lambda(t) = \begin{bmatrix} \lambda_{11}(t) & \dots & \lambda_{1n}(t) \\ \vdots & & \vdots \\ \lambda_{n1}(t) & \dots & \lambda_{nn}(t) \end{bmatrix}$$

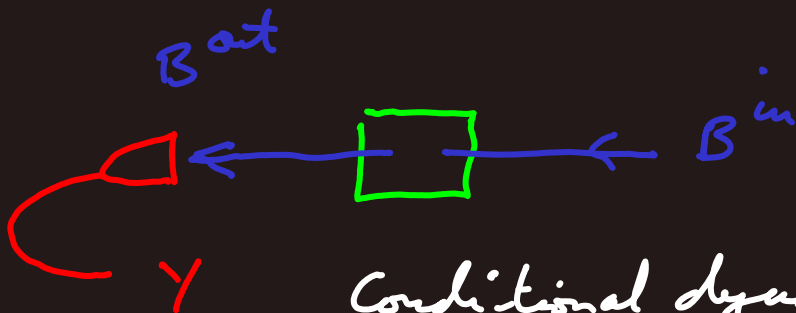
$$S = \begin{bmatrix} s_{11} & \dots & s_{1n} \\ \vdots & & \vdots \\ s_{n1} & \dots & s_{nn} \end{bmatrix}, \quad L = \begin{bmatrix} L_1 \\ \vdots \\ L_n \end{bmatrix}, \quad H = H^*$$

$G \sim (S, L, H)$

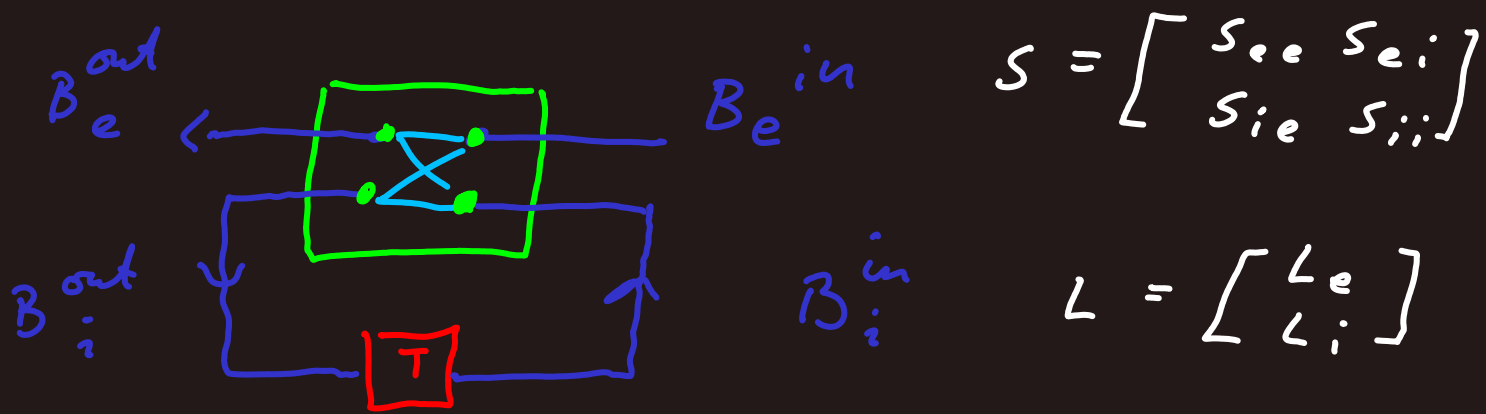
Input - Output.



$$B_k^{out}(t) = U_t^* (I \otimes B_n(t)) U_t$$



Conditional dynamics of the system due to output measurement



$$\begin{cases} S_{fb} = S_{ee} + S_{ei} T (1 - s_{ii} T)^{-1} S_{ie} \\ L_{fb} = L_e + S_{ei} T (1 - s_{ii} T)^{-1} L_i \\ H_{fb} = H + \sum_{\alpha=e,i} I_m L_{\alpha}^* S_{\alpha i} T (1 - s_{ii} T)^{-1} L_i \end{cases}$$

→ Complete Foundation for open quantum input-output models

- * interconnection
- * feedback
- * conditional evolution (filtering)
- * Control / System Ident.